SUSY CP Problem in Gauge Mediation Model

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Supersymmetry

- SUSY is motivated by facts that SUSY ...
 - Solves the hierarchy problem
 - Preferred by gauge coupling unification
 - Provide a candidate for dark matter with R-parity

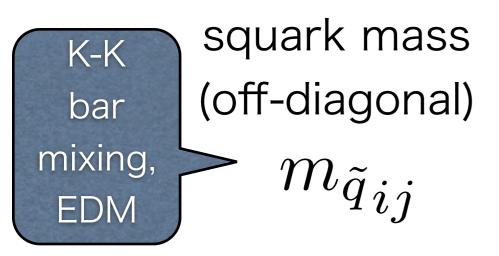
However ...

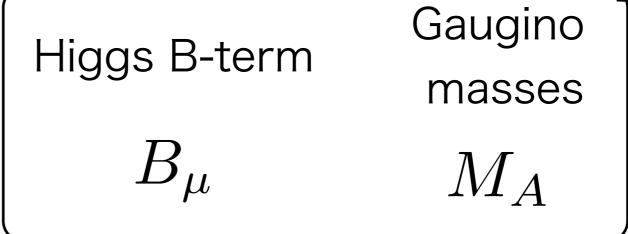
SUSY CP Problem

In terms of the CP violation, soft SUSY breaking terms can cause trouble

Dangerous sources for CP violation

this talk

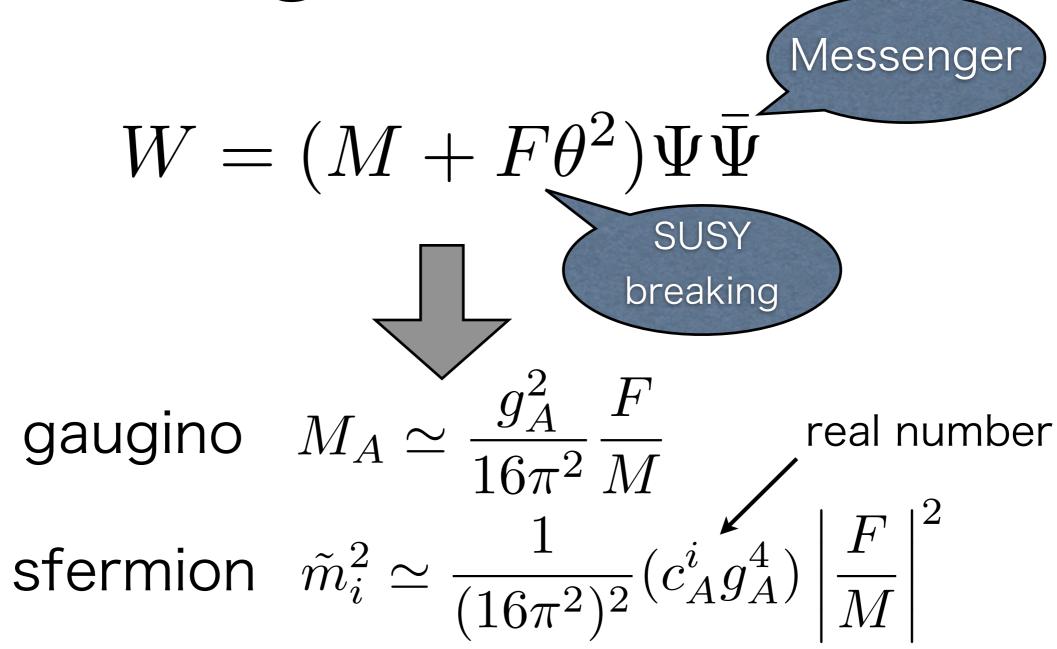




(The phases of A-terms are also relevant)

How to generate soft SUSY breaking parameters is important

Gauge Mediation



Both F and M can be taken to be real

No CPV in gaugino and sfermion masses

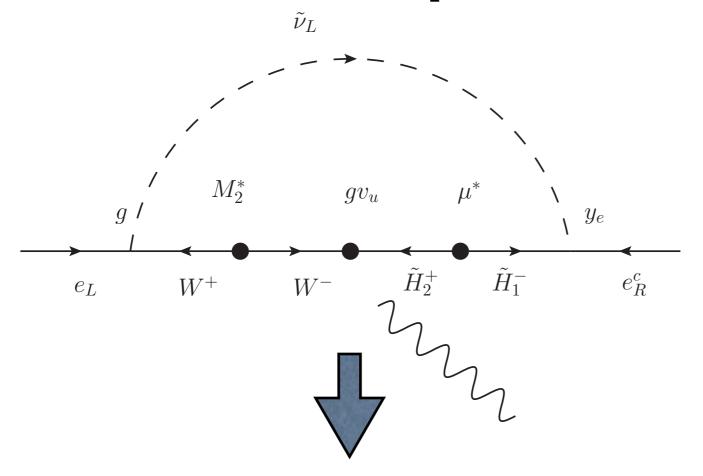
However, even in Gauge Mediation, new CPV phase can be induced

2. Small GUT breaking operator (Plank suppressed)

3. SUGRA effects

In this talk, we focus on 2 and 3

Electric Dipole Moment



chargino diagram is dominant

(unless higgsino and wino are heavy)

 $d_e \propto m_e \tan \beta \operatorname{Arg}(M_2 B_{\mu}^*)/m_{soft}^2$

proportional to relative phase between Wino mass and Higgs B-term

exp. constraint

 $d_e < 2.1 \times 10^{-27} e \,\mathrm{cm}$ (95% C.L.)

The phase should be small as $O(10^{-3}-10^{-4})$ with O(1) TeV sparticles and O(10) tan β

Muon g-2

When sparticles are heavy the constraint on the phase becomes loose, however

The EXP value and SM prediction of muon g-2 are deviated more than 3σ level

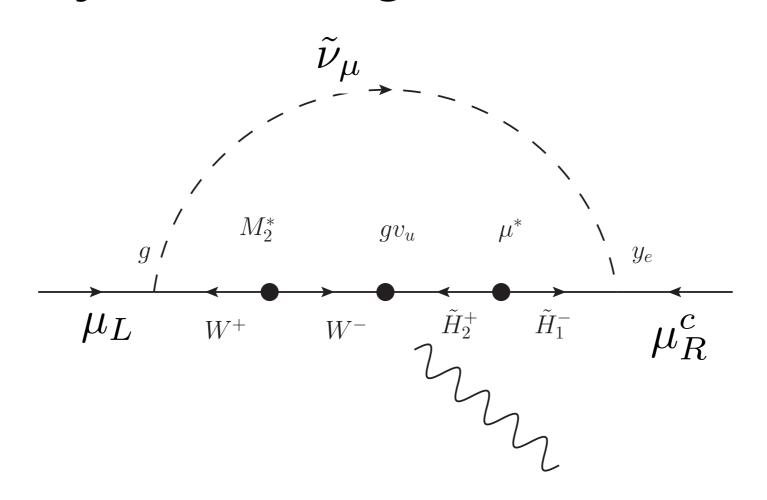
$$a_{\mu}^{\rm EXP} - a_{\mu}^{\rm SM} = (26.1 \pm 8.0) \cdot 10^{-10}$$
 3.3 σ deviation

(Hagiwara, Martin, Nomura, Teubner, 2011)

The other group also reported the discrepancy more than 3σ level (M. Davier, A. Hoecker, B. Malaescu, Z. Zhang, 2010)

If the SUSY is responsible for the deviation of muon g-2, sleptons can not be heavy

Dominant SUSY contribution to the muon g-2 is essentially same diagram as that of EDM



$$d_e \simeq rac{m_e}{2m_\mu^2} {
m Arg}(M_2 B_\mu^*) a_\mu^{
m SUSY}$$
 (J. L. Feng, T. Moroi, 1999)

To suppress the EDM while explaining the muon g-2 is quite difficult

 $\left(\operatorname{Arg}(B_{\mu}) \right)$

 μ/B_{μ} have to be induced with the phase smaller than O(10⁻³ - 10⁻⁴)

(e.g.) Only SUSY mass, μ is generated above the messenger scale

$$W=\lambda \frac{X^2}{M_P} H_u H_d$$
 PQ breaking scale $10^9 {\rm GeV} \lesssim \langle X \rangle \lesssim 10^{12} {\rm GeV}$

However, $O(10^{-3} - 10^{-4})$ phase can arise in other ways

The CPV effects from Dim. 5 GUT breaking operator

24 Higgs case

$$\langle \Sigma \rangle$$

 $SU(5) \to SU(3)_C \times SU(2)_L \times U(1)_Y$

24 Higgs breaks GUT symmetry

$$\langle \Sigma \rangle = \frac{v_{24}}{2\sqrt{15}} \operatorname{diag}(2, 2, 2, -3, -3)$$

Messenger

(The GUT breaking SUSY mass term may also exist)

Messenger dim. 5
$$W = \lambda_0 S \bar{\Psi}^{\alpha} \Psi_{\alpha} + \frac{\lambda_1}{M_{\rm Pl}} S \bar{\Psi}^{\alpha} \Sigma_{\alpha}^{\beta} \Psi_{\beta} + M \bar{\Psi}^{\alpha} \Psi_{\alpha}$$

$$= \lambda_0 \left(1 + \frac{1}{\sqrt{15}} \epsilon_{24} \right) S \bar{\psi}_d \psi_d + \lambda_0 \left(1 - \frac{3}{2\sqrt{15}} \epsilon_{24} \right) S \bar{\psi}_l \psi_l + \cdots$$

complex

$$\epsilon_{24} = \frac{\lambda_1 v_{24}}{\lambda_0 M_{\rm Pl}} \sim \mathcal{O}(10^{-2})$$

The messengers of SU(2) doublet and SU(3) triplet have different phases of O(10⁻³)

The phases of gauginos are not aligned any more

$$M_{1} = \frac{g_{1}^{2}}{16\pi^{2}} \left(1 - \frac{1}{2\sqrt{15}} \epsilon_{24} \right) \frac{F_{S}}{M}$$

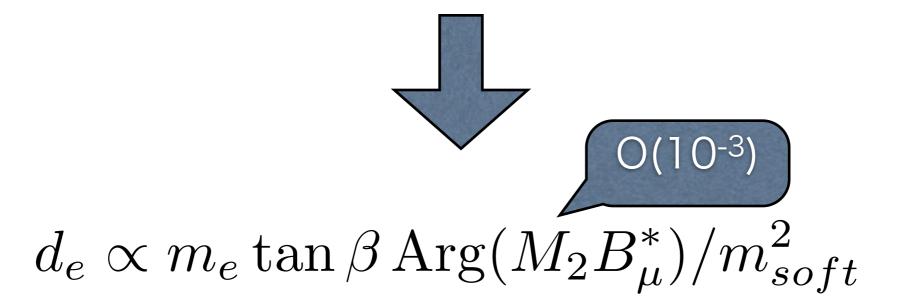
$$M_{2} = \frac{g_{2}^{2}}{16\pi^{2}} \left(1 - \frac{3}{2\sqrt{15}} \epsilon_{24} \right) \frac{F_{S}}{M} \qquad \begin{array}{c} \text{complex} \\ \epsilon_{24} = \frac{\lambda_{1} v_{24}}{\lambda_{0} M_{\text{Pl}}} \sim \mathcal{O}(10^{-2}) \end{array}$$

$$M_{3} = \frac{g_{2}^{2}}{16\pi^{2}} \left(1 + \frac{1}{\sqrt{15}} \epsilon_{24} \right) \frac{F_{S}}{M}$$

The phases differ by O(10⁻³)

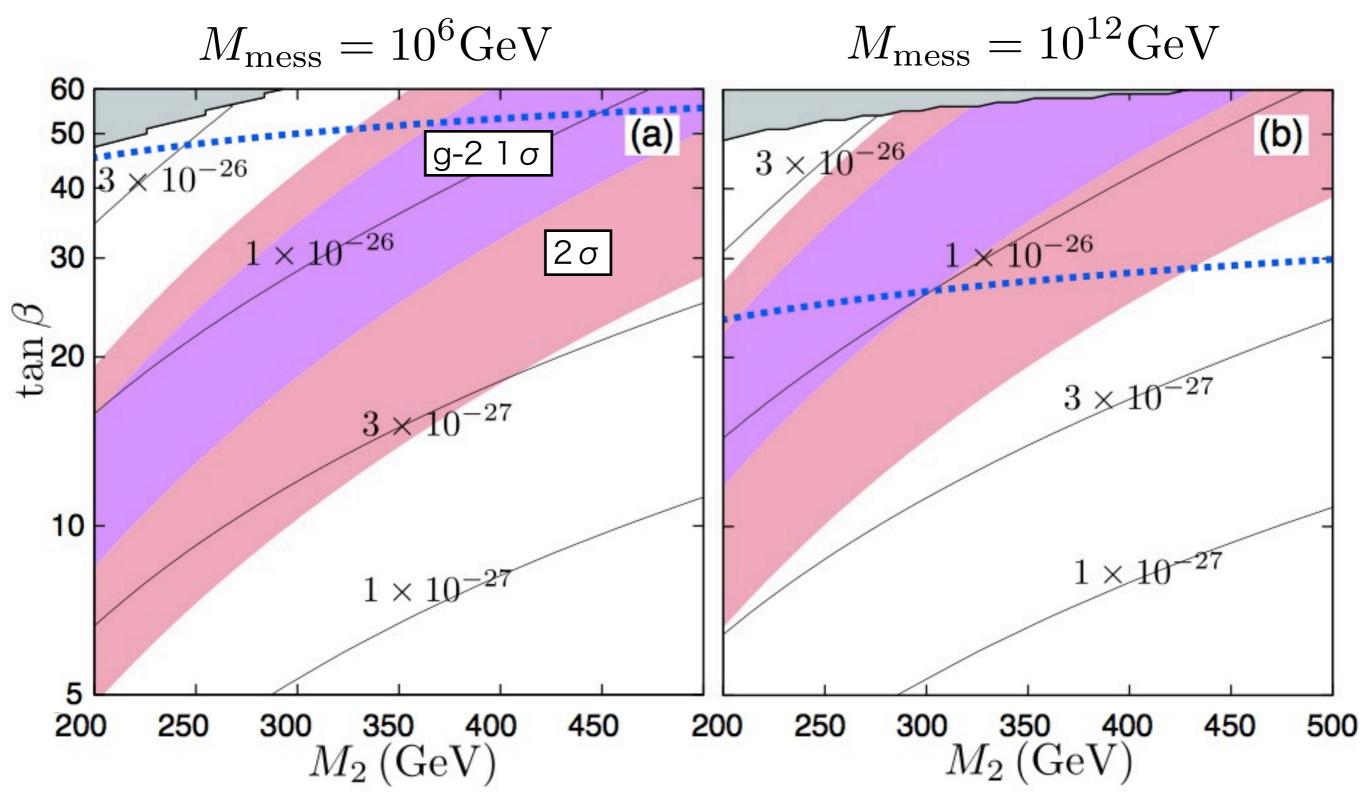
The Higgs B-term is also affected by the phase difference of gauginos through RGE

$$\frac{dB_{\mu}}{d\ln\mu} = \frac{1}{8\pi^2} \left[3g_2^2 M_2 + g_1^2 M_1 - 3\text{tr}(Y_U^{\dagger} A_U) - 3\text{tr}(Y_D^{\dagger} A_D) - \text{tr}(Y_L^{\dagger} A_L) \right]$$



The induced CP phase is large enough to be constrained from EDM experiments

for numerical calculation $\epsilon_{24} = i0.01$



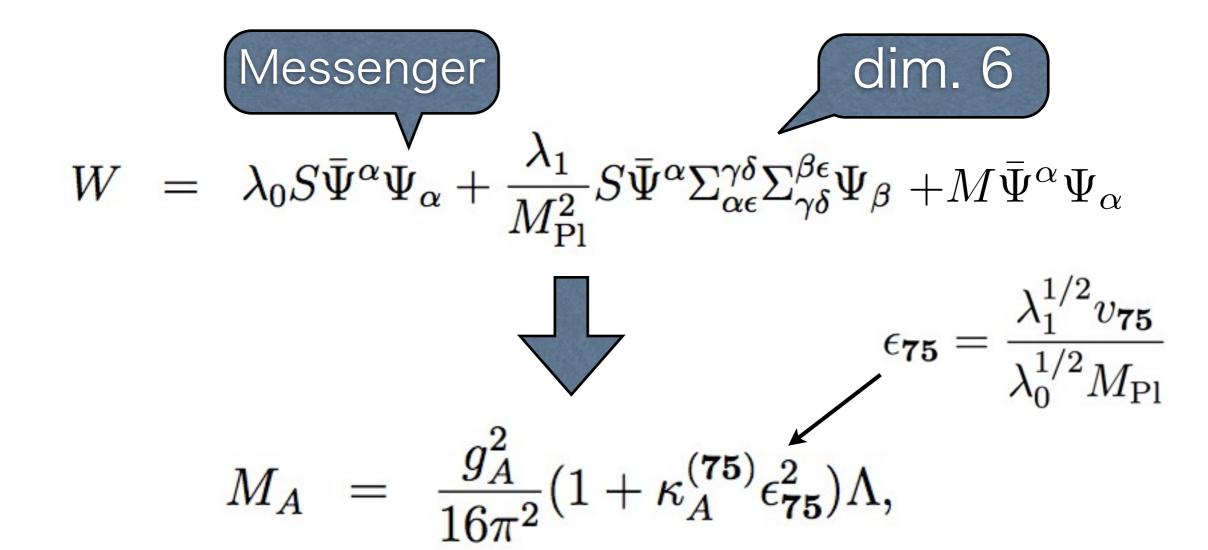
The dim 5 GUT breaking operator should be suppressed by O(0.1) somehow exp. bound $d_e < 2.1 \times 10^{-27} e~{\rm cm}$ (95% C.L.).

The CPV effects from Dim. 6 GUT breaking operator

75 Higgs case

$$\langle \Sigma \rangle$$

 $SU(5) \to SU(3)_C \times SU(2)_L \times U(1)_Y$



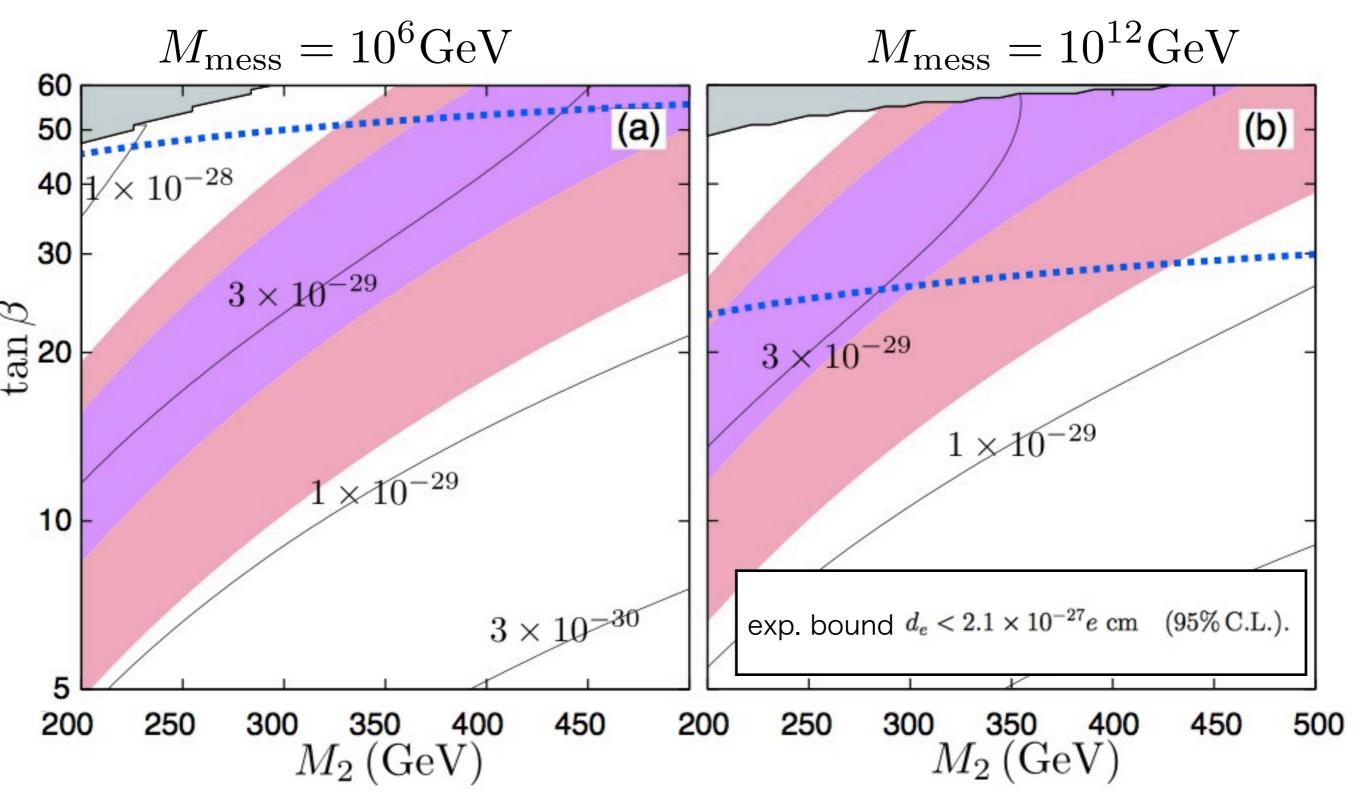
$$\kappa_1^{(75)} = \frac{11}{60}, \, \kappa_2^{(75)} = \frac{1}{4}, \, \kappa_3^{(75)} = \frac{1}{12}.$$

The induced CP phase

 $O(10^{-5})$

 $d_e \propto m_e \tan \beta \operatorname{Arg}(M_2 B_{\mu}^*)/m_{soft}^2$

for numerical calculation $\epsilon_{75} = i0.01$



The dim 6 GUT breaking operator is consistent with g-2 and EDM
EDM may be seen at the future experiment

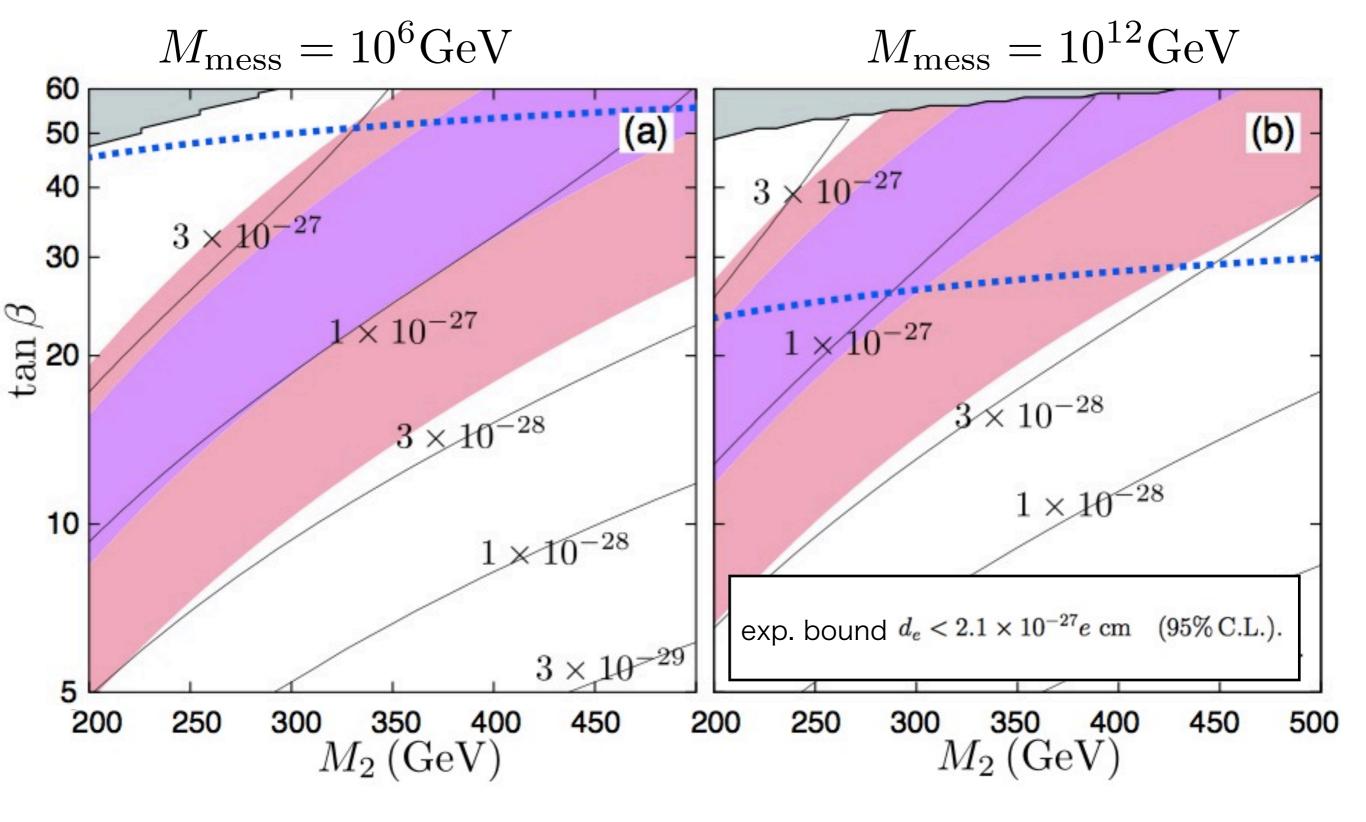
The SUGRA effects

- Even in GMSB, the effects from SUGRA exists
- Higgs B-term obtains additional contribution from SUGRA

$$B_{\mu}=B_{\mu}^{(0)}+B_{\mu}^{(\mathrm{SUGRA})}$$
 complex
$$B_{\mu}^{(\mathrm{SUGRA})}\sim m_{3/2}$$

For numerical calculation, we take

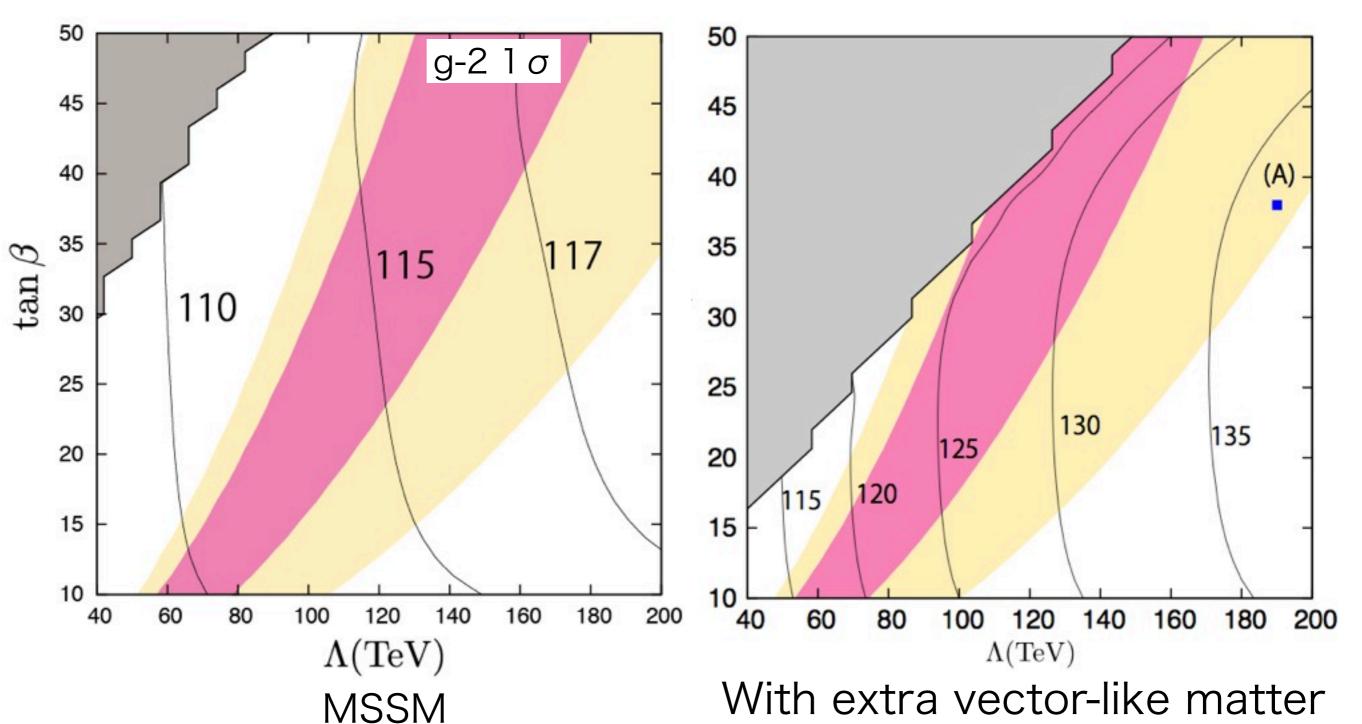
$$B_{\mu}^{(SUGRA)} = i100 \text{MeV}$$



100 MeV gravitino is marginally consistent with g-2 and EDM

Higgs mass in minimal GMSB

 $M_{
m mess}=5 imes10^5{
m GeV}$ and $N_5=1$



(M. Endo, K. Hamaguchi, S. Iwamoto, N. Yokozaki, 2011)

Summary

- The effects of the GUT breaking operator is sizable, and may induce large EDM
- The effect of gravity mediation is also sizable when the gravitino mass is large as 100MeV
- EDM may be seen at future experiments